

Potential in a polygonal body coated  
with a thin dielectric layer

Grégory VIAL

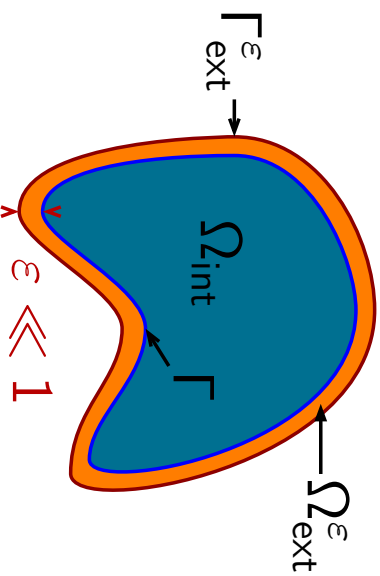
IRMAR, Rennes, France.

Oberwolfach, November 2002.

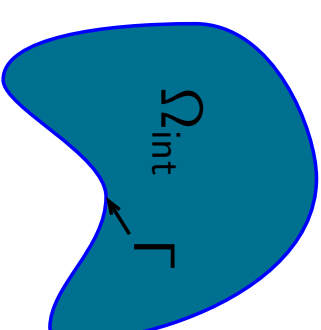
## Overview

- Approximate boundary conditions : the smooth case.
- Case of a corner domain : what to do with singularities.
- Numerical experiments.

$$\left\{ \begin{array}{lll} \alpha \Delta u_{\text{int}}^\varepsilon & = & -f_{\text{int}} \quad \text{in } \Omega_{\text{int}}, \\ \Delta u_{\text{ext}}^\varepsilon & = & 0 \quad \text{in } \Omega_{\text{ext}}^\varepsilon, \\ u_{\text{int}}^\varepsilon & = & u_{\text{ext}}^\varepsilon \quad \text{on } \Gamma, \\ \alpha \partial_n u_{\text{int}}^\varepsilon & = & \partial_n u_{\text{ext}}^\varepsilon \quad \text{on } \Gamma, \\ u_{\text{ext}}^\varepsilon & = & 0 \quad \text{on } \Gamma_{\text{ext}}^\varepsilon. \end{array} \right.$$



$$\left\{ \begin{array}{ll} \alpha \Delta v^\varepsilon = -f_{\text{int}} & \text{in } \Omega_{\text{int}}, \\ \text{BC}(\varepsilon) & \text{on } \Gamma. \end{array} \right.$$



Choice of  $\text{BC}(\varepsilon)$

such that  $u_{\text{int}}^\varepsilon - v^\varepsilon = \mathcal{O}(\varepsilon^p)$  ?

p 1

2

3

$$\text{BC}(\varepsilon) \quad v^\varepsilon = 0 \quad v^\varepsilon + \varepsilon \alpha \partial_n v^\varepsilon = 0 \quad \left( 1 + \varepsilon \frac{c(t)}{2} \right) v^\varepsilon + \varepsilon \alpha \partial_n v^\varepsilon = 0$$

c(t) is the curvature on  $\Gamma$  on the point of curvilinear abscissa t.

We perform a dilatation of  $\Omega_{\text{ext}}^\varepsilon$  in the normal direction :

$$\left. \begin{array}{ccc} \Omega_{\text{ext}}^\varepsilon & \xrightarrow{\quad} & \Omega_{\text{ext}}^1 \\ (s, t) & \longmapsto & (\bar{s} = \frac{s}{\varepsilon}, t) \end{array} \right\} (s, t : \text{Frénet variables}).$$

We define :  $\bar{u}_{\text{ext}}^\varepsilon(\bar{s}, t) = u_{\text{ext}}^\varepsilon(s, t)$ .

We consider a fixed domain ;  $\varepsilon$  appears now in the equations.

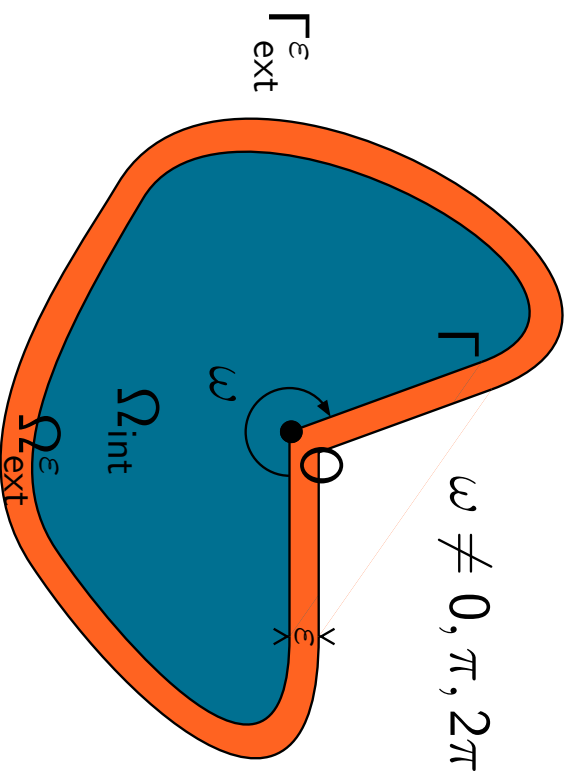
$$\left\{ \begin{array}{ll} \alpha \Delta u_{\text{int}}^\varepsilon = -f_{\text{int}} & \text{in } \Omega_{\text{int}}, \\ \frac{1}{\varepsilon^2} \left( \partial_s^2 - \sum_{\ell \geq 1} \varepsilon^\ell A^\ell \right) \bar{u}_{\text{ext}}^\varepsilon = 0 & \text{in } \Omega_{\text{ext}}^1, \\ u_{\text{int}}^\varepsilon = \bar{u}_{\text{ext}}^\varepsilon & \text{on } \Gamma, \\ \alpha \partial_n u_{\text{int}}^\varepsilon = \frac{1}{\varepsilon} \partial_s \bar{u}_{\text{ext}}^\varepsilon & \text{on } \Gamma, \\ \bar{u}_{\text{ext}}^\varepsilon = 0 & \text{on } \Gamma_{\text{ext}}^1. \end{array} \right.$$

$$\left\{ \begin{array}{ll} \alpha \Delta u_{\text{int}}^\varepsilon = -f_{\text{int}} & \text{in } \Omega_{\text{int}}, \\ \frac{1}{\varepsilon^2} \left( \partial_s^2 - \sum_{\ell \geq 1} \varepsilon^\ell A^\ell \right) \bar{u}_{\text{ext}}^\varepsilon = 0 & \text{in } \Omega_{\text{ext}}^1, \\ u_{\text{int}}^\varepsilon = \bar{u}_{\text{ext}}^\varepsilon & \text{on } \Gamma, \\ \alpha \partial_n u_{\text{int}}^\varepsilon = \frac{1}{\varepsilon} \partial_s \bar{u}_{\text{ext}}^\varepsilon & \text{on } \Gamma, \\ \bar{u}_{\text{ext}}^\varepsilon = 0 & \text{on } \Gamma_{\text{ext}}^1. \end{array} \right.$$

The ansätze  $u_{\text{int}}^\varepsilon = \sum_{n \geq 0} \varepsilon^n u_{\text{int}}^n$  and  $\bar{u}_{\text{ext}}^\varepsilon = \sum_{n \geq 0} \varepsilon^n \bar{u}_{\text{ext}}^n$  lead to the problems

$$\left\{ \begin{array}{ll} \partial_s^2 \bar{u}_{\text{ext}}^n = \sum_{\ell+p=n} A^\ell \bar{u}_{\text{ext}}^{-p} & \text{in } \Omega_{\text{ext}}^1, \\ \partial_s \bar{u}_{\text{ext}}^n = \alpha \partial_n u_{\text{int}}^{n-1} & \text{on } \Gamma, \\ \bar{u}_{\text{ext}}^n = 0 & \text{on } \Gamma_{\text{ext}}^1. \end{array} \right. \quad \left\{ \begin{array}{ll} \alpha \Delta u_{\text{int}}^n = -f_{\text{int}} \delta_0^n & \text{in } \Omega_{\text{int}}, \\ u_{\text{int}}^n = \bar{u}_{\text{ext}}^n & \text{on } \Gamma. \end{array} \right.$$

We obtain the asymptotic expansion of  $u^\varepsilon$ ; by truncature of the series we deduce the approximate boundary conditions.



We perform a dilatation of  $\Omega_{\text{ext}}^\epsilon$  in the “normal” direction.

Ansätze :  $\mathbf{u}_{\text{int}}^\epsilon = \sum_{n \geq 0} \epsilon^n \mathbf{u}_{\text{int}}^n$

$$\bar{\mathbf{u}}_{\text{ext}}^\epsilon = \sum_{n \geq 0} \epsilon^n \bar{\mathbf{u}}_{\text{ext}}^n$$

$$\left\{ \begin{array}{l} \partial_s^2 \bar{\mathbf{u}}_{\text{ext}}^n = \sum_{\ell+p=n} A^\ell \bar{\mathbf{u}}_{\text{ext}}^p \quad \text{in } \Omega_{\text{ext}}^1, \\ \partial_s \bar{\mathbf{u}}_{\text{ext}}^n = \alpha \partial_n \mathbf{u}_{\text{int}}^{n-1} + \mathbf{a}(t) \partial_t \bar{\mathbf{u}}_{\text{ext}}^{n-1} \quad \text{on } \Gamma, \\ \bar{\mathbf{u}}_{\text{ext}}^n = 0 \quad \text{on } \Gamma_{\text{ext}}^1. \end{array} \right.$$

$$\left\{ \begin{array}{l} \alpha \Delta \mathbf{u}_{\text{int}}^n = -\mathbf{f}_{\text{int}} \delta_0^n \quad \text{in } \Omega_{\text{int}}, \\ \mathbf{u}_{\text{int}}^n = \bar{\mathbf{u}}_{\text{ext}}^n \quad \text{on } \Gamma. \end{array} \right.$$

- At each step, we loose some regularity.
- Singularities appear, due to the corner.

For  $n = 0$ 

$$\bar{u}_{\text{ext}}^0 = 0 \quad \text{and}$$

$$\begin{cases} \alpha \Delta u_{\text{int}}^0 = -f_{\text{int}} & \text{in } \Omega_{\text{int}}, \\ u_{\text{int}}^0 = 0 & \text{on } \Gamma. \end{cases}$$

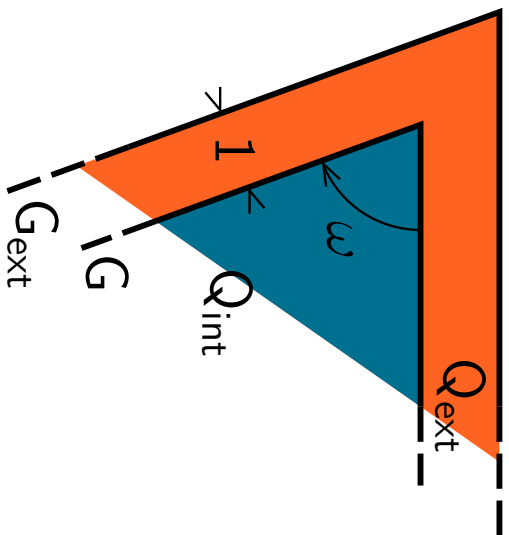
The regularity of  $u_{\text{int}}^0$  is limited by the first singular function :  $\chi s^1(r, \theta)$  ;

- $\chi$  : cut-off function equal to 1 near the corner ;
- $s^1(r, \theta) = r^{\frac{\pi}{\omega}} \sin\left(\frac{\pi}{\omega}\theta\right)$  ,  $(r, \theta)$  : polar coordinates around the corner.

For  $n = 1$ 

$$\bar{u}_{\text{ext}}^1 = \alpha \partial_n u_{\text{int}}^0(\bar{s} - 1) \notin H^1(\Omega_{\text{ext}}^1) \quad \text{if } \omega \geq \frac{2\pi}{3}.$$

*Due to the loss of regularity, the singularities cannot be handled with the classical method (exterior-interior successive resolutions).*



- $u_{\text{int}}^0 \text{ flat} \implies u_{\text{int}}^0 = u_{\text{flat,int}}^0 + \chi \sum_q c_q^0 s^q(r, \theta)$  (\*)

where  $s^q(r, \theta) = r^{\frac{q\pi}{\omega}} \sin\left(\frac{q\pi}{\omega}\theta\right)$ ,  $q \geq 1$ .

- $s^q$  solves the homogeneous Dirichlet problem in  $Q_{\text{int}}$ .
- Let  $\mathcal{R}^q$  be solution of the homogeneous transmission problem in  $Q_{\text{int}}$  and  $Q_{\text{ext}}$  such that  $\mathcal{R}^q \underset{r \rightarrow +\infty}{\sim} s^q$ .

In (\*), we replace the singularities

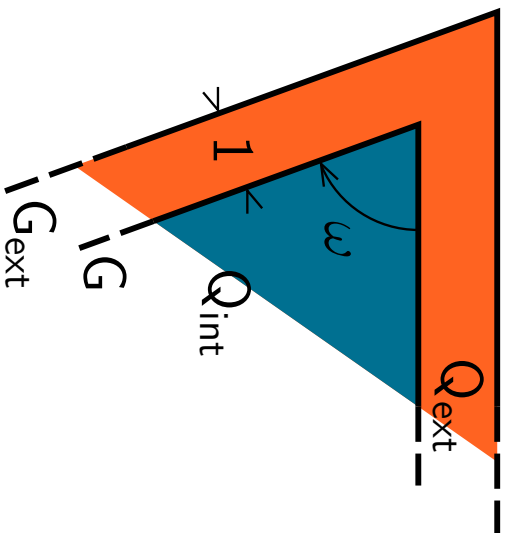
$$u_{\text{int}}^0 = u_{\text{flat,int}}^0 + \chi \sum_q c_q^0 \varepsilon^{\frac{q\pi}{\omega}} s^q\left(\frac{r}{\varepsilon}, \theta\right) \text{ becomes } u_{\text{flat,int}}^0 + \chi \sum_q c_q^0 \varepsilon^{\frac{q\pi}{\omega}} \mathcal{R}^q\left(\frac{r}{\varepsilon}, \theta\right)$$

The part  $u_{\text{flat,int}}^0$  is handled with the “exterior-interior” method.

The final asymptotic expansion is

$$u^\varepsilon \approx \sum_\mu \varepsilon^\mu u_{\text{flat}}^\mu + \chi \sum_q \sum_\mu \varepsilon^{\mu + \frac{q\pi}{\omega}} c_q^\mu \mathcal{R}^q\left(\frac{r}{\varepsilon}, \theta\right), \mu = p + \frac{k\pi}{\omega} \quad \left| \begin{array}{l} p \in \mathbf{N} \\ k = 0 \text{ or } k \geq 2 \end{array} \right.$$

First exponents :  $\mu = 0, 1, \frac{2\pi}{\omega}, \frac{3\pi}{\omega}, 2, \dots$



Variational space

$$\left\{ \begin{array}{ll} \alpha \Delta \mathcal{R}_{\text{int}}^q = 0 & \text{in } Q_{\text{int}}, \\ \Delta \mathcal{R}_{\text{ext}}^q = 0 & \text{in } Q_{\text{ext}}, \\ \mathcal{R}_{\text{int}}^q = \mathcal{R}_{\text{ext}}^q & \text{on } G, \\ \alpha \partial_n \mathcal{R}_{\text{int}}^q = \partial_n \mathcal{R}_{\text{ext}}^q & \text{on } G, \\ \mathcal{R}_{\text{ext}}^q = 0 & \text{on } G_{\text{ext}}, \\ \mathcal{R}^q \sim \mathfrak{s}_0^q & \text{when } r \rightarrow \infty. \end{array} \right.$$

$$\mathfrak{X} = \left\{ \mathbf{v} ; \frac{\mathbf{v}}{1+r}, \nabla \mathbf{v} \in L^2(Q) \text{ and } \mathbf{v}_{\text{ext}} = 0 \text{ on } G_{\text{ext}} \right\} .$$

Problem

if  $q \geq 0$ ,  $\mathfrak{s}_0^q$  does not belong to  $\mathfrak{X}$  !

### Over-variational expansion

We construct  $\mathfrak{R}^{(0)}, \dots, \mathfrak{R}^{(p)}$  ( $p \geq \frac{q\pi}{\omega}$ ) such that there exists  $\mathfrak{v} \in \mathfrak{M}$  and

$$\mathfrak{R}^q := \mathfrak{v} + \mathfrak{R}^{(0)} + \dots + \mathfrak{R}^{(p)}$$

satisfies the problem.

### Under-variational expansion

The Mellin transform  $\hat{v}(\lambda)$  of  $\mathfrak{v}$  is defined for  $\text{Re } \lambda > 0$ .

- $\hat{v}(\lambda)$  has a meromorphic extension  $\kappa$  in  $\mathbb{C}$  ;
- If  $\mathcal{M}_\xi^{-1}$  denotes the inverse Mellin transformation along the line  $\text{Re } \lambda = \xi$ ,

$$\mathfrak{v} - \mathcal{M}_\xi^{-1}(r^\lambda \kappa) = \sum \text{Res } r^\lambda \kappa(\lambda).$$

### Complete expansion

There exists a solution  $\mathfrak{R}^q$  to the problem in the infinite domain such that :

$$\mathfrak{R}^q = \mathfrak{R}^{(0)} + \dots + \mathfrak{R}^{(p)} + \mathfrak{R}^{(p+1)} + \dots + \mathfrak{R}^{(p+N)} + o_\infty\left(r^{\frac{q\pi}{\omega} - p - N}\right).$$

For the transmission problem

$$u_{\text{int}}^\varepsilon \approx \sum_{\mu} \varepsilon^\mu u_{\text{flat,int}}^\mu + \chi \sum_q \sum_{\mu} c_q^\mu \varepsilon^{\mu + \frac{q\pi}{\omega}} \mathcal{R}^q\left(\frac{r}{\varepsilon}, \theta\right)$$

For the impedance problem ( $v^\varepsilon + \varepsilon \alpha \partial_n v^\varepsilon = 0$ )

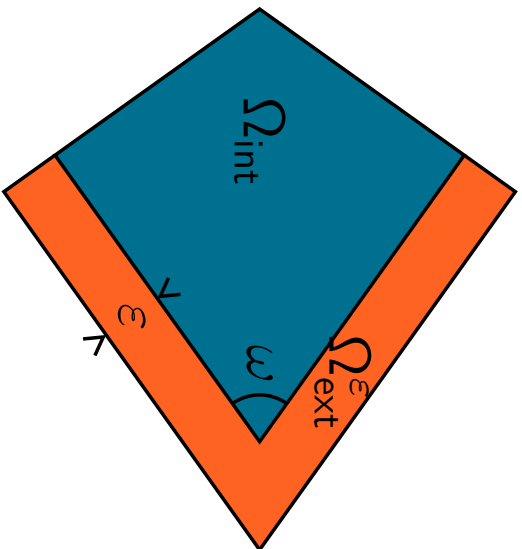
$$v^\varepsilon \approx \sum_{\mu} \varepsilon^\mu v_{\text{flat}}^\mu + \chi \sum_q \sum_{\mu} d_q^\mu \varepsilon^{\mu + \frac{q\pi}{\omega}} \mathcal{I}^q\left(\frac{r}{\varepsilon}, \theta\right)$$

- For  $\mu < \min(3, 1 + \frac{2\pi}{\omega})$ ,  $u_{\text{flat,int}}^\mu = v_{\text{flat}}^\mu$  and  $c_q^\mu = d_q^\mu$ .
- The profiles satisfy  $\mathcal{R}^q - \mathcal{I}^q = \mathcal{O}\left(r^{\max(\frac{q\pi}{\omega} - 3, -\frac{\pi}{\omega})}\right)$  as  $r \rightarrow +\infty$ .

### Error estimates

$$\|u_{\text{int}}^\varepsilon - v^\varepsilon\|_{L^2(\Omega_{\text{int}})} = \mathcal{O}\left(\varepsilon^{\min(\frac{2\pi}{\omega}, 1 + \frac{\pi}{\omega}, 3)}\right)$$

$$\|u_{\text{int}}^\varepsilon - v^\varepsilon\|_{H^1(\Omega_{\text{int}})} = \mathcal{O}\left(\varepsilon^{\min(\frac{\pi}{\omega}, 3)}\right)$$



We compute :

- the solution of the transmission problem,
- the solution of the impedance problem.

$\alpha = 10$  and  $\text{supp } f_{\text{int}} \subset \subset \Omega_{\text{int}}$ .

15 values of  $\varepsilon$  and 12 values of  $\omega$

Computations done with the f.e. library **MÉLINA** [Daniel Martin, IRMAR]

Using

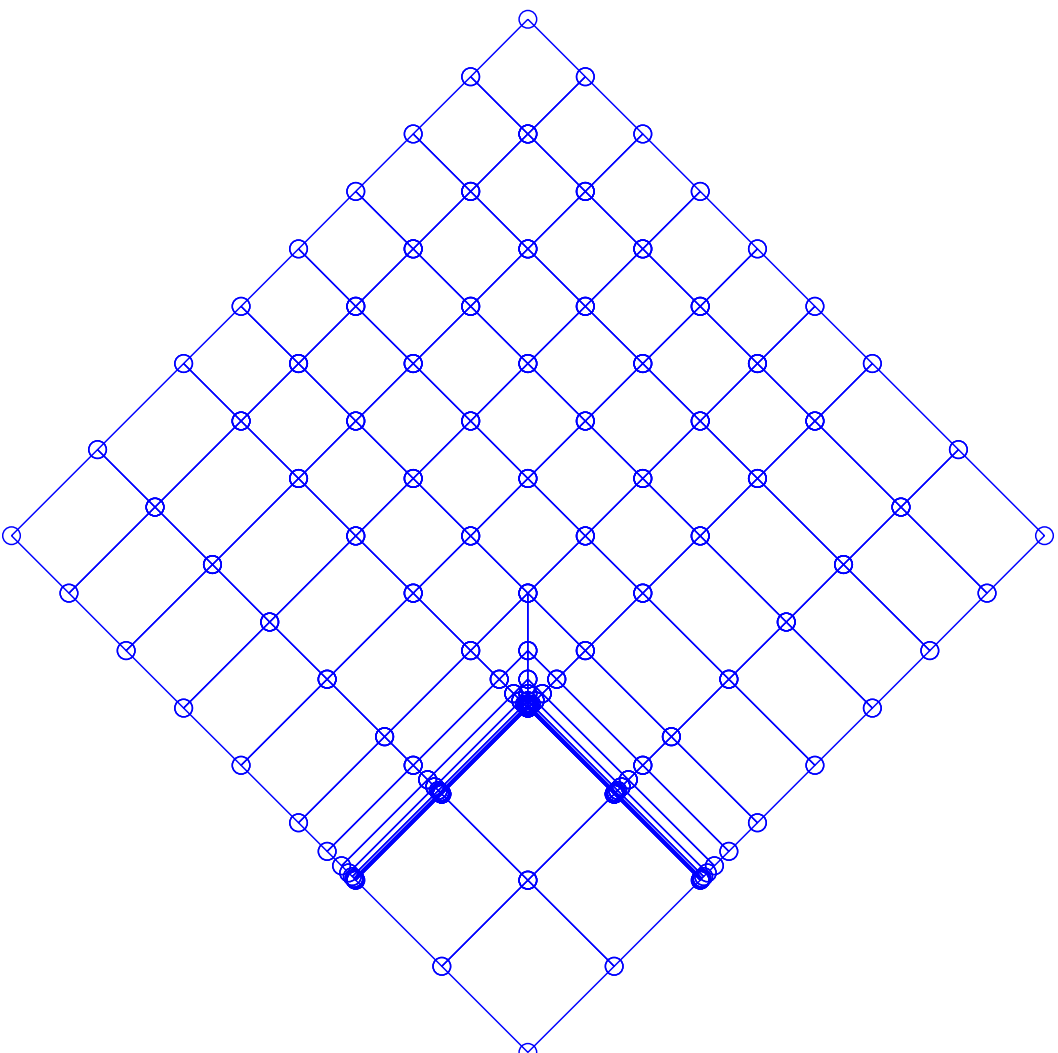
Lagrange  $\mathbb{Q}_8$  finite elements,

Geometric refinement near the corner.

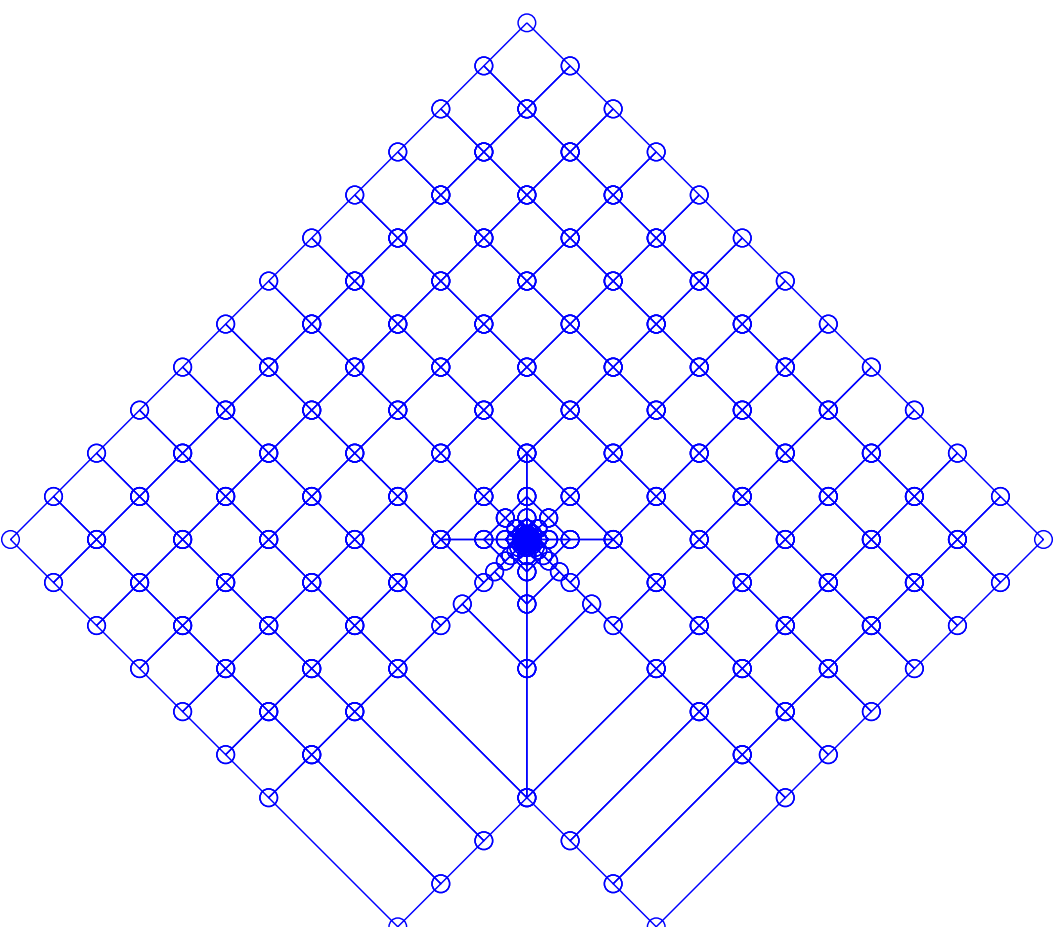
allows

Accuracy of the computation,

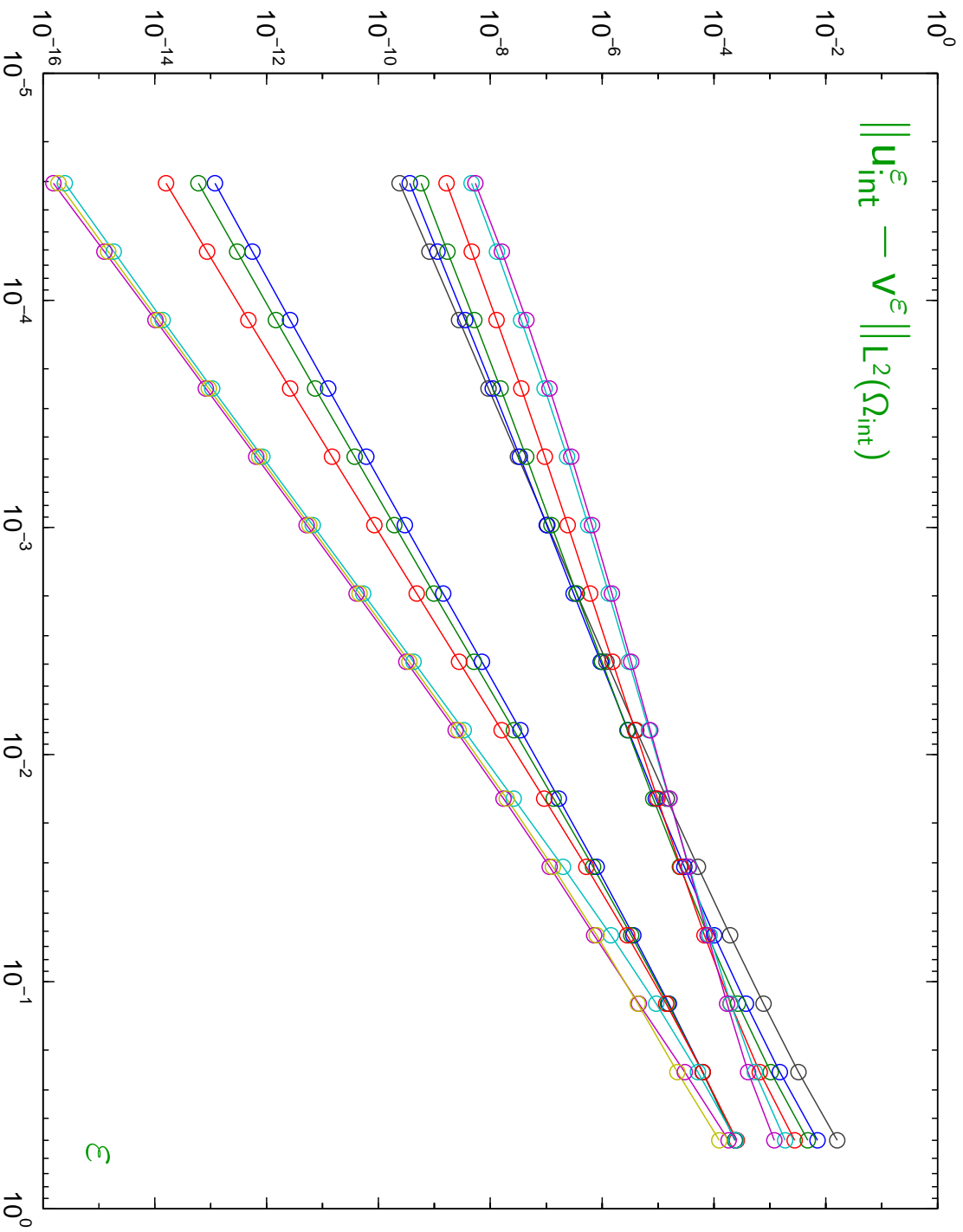
Possible anisotropy in the layer.



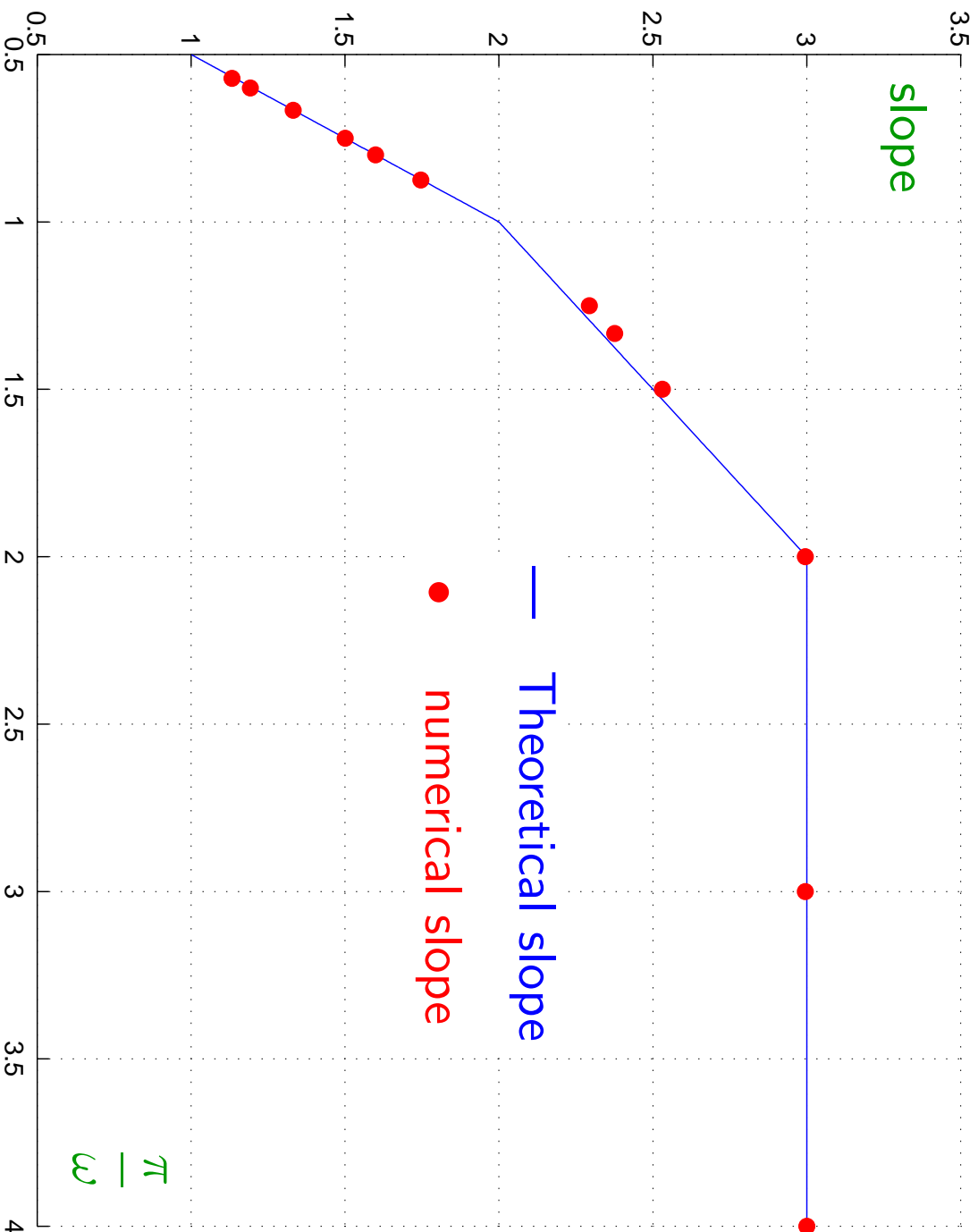
Mesh for  $\omega = \frac{\pi}{2}$



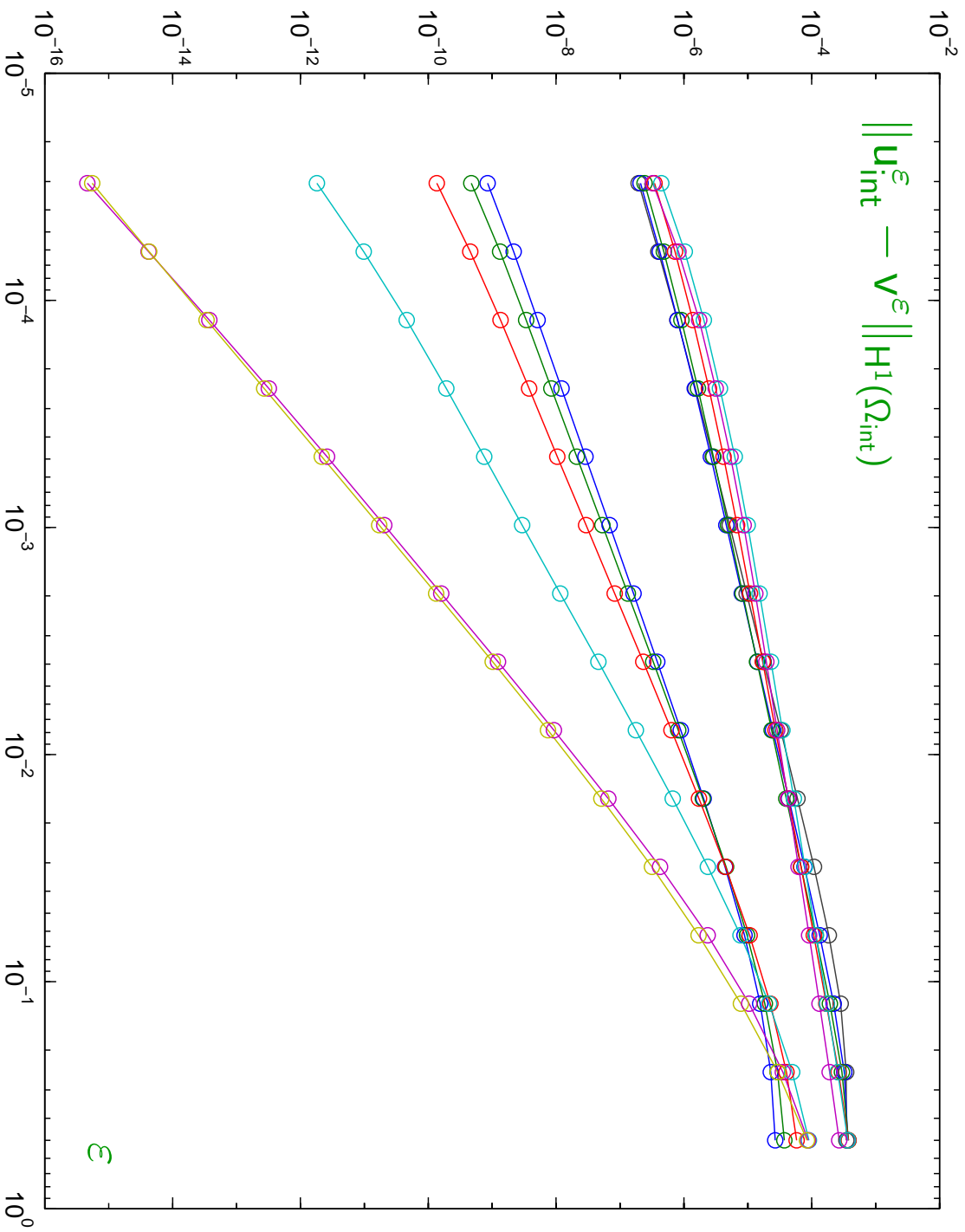
Mesh for  $\omega = \frac{3\pi}{2}$



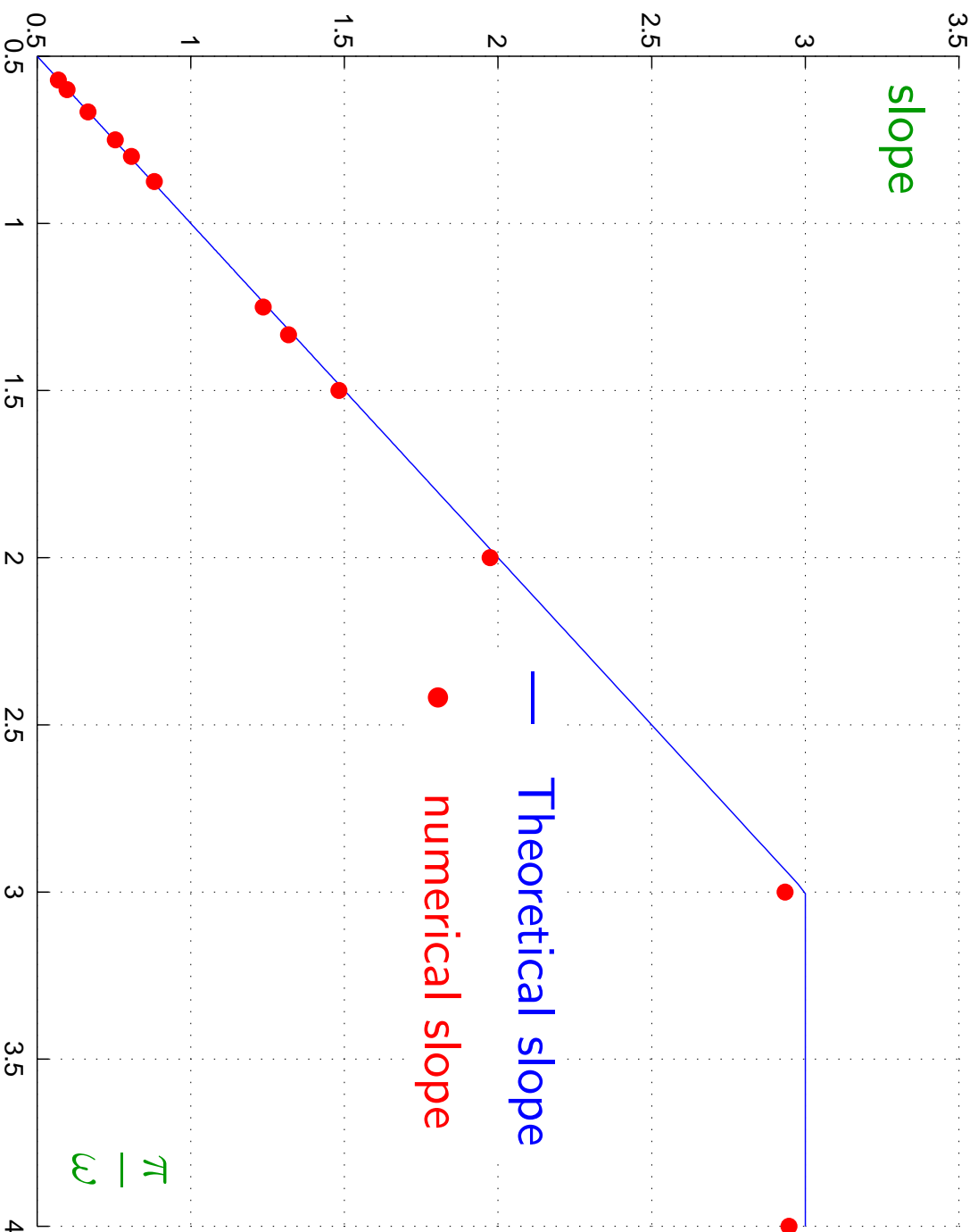
Each “line” corresponds to a value of  $\omega$ .



Theoretical and numeric slopes ( $L^2$  norm)



Each “line” corresponds to a value of  $\omega$ .



Theoretical and numeric slopes ( $H^1$  norm)